## TWO MODEL HAMILTONIANS FOR TORSION-INVERSION TUNNELING IN THE CH-STRETCH VIBRA-TIONALLY EXCITED STATES OF METHYLAMINE

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In methylamine (CH<sub>3</sub>NH<sub>2</sub>), there are six equivalent mimina that are connected by torsion and inversion tunneling. In the  $G_{12}$  molecular symmetric group, there are four species,  $A = \{A_1, A_2\}$ ,  $B = \{B_1, B_2\}$ ,  $E_1$  and  $E_2$  that combine with distinct nuclear states. The ground vibrational state of CH<sub>3</sub>NH<sub>2</sub> is split by torsion and inversion tunneling into a multiplet pattern of four distinct energy levels<sup>*a*</sup>. The experimental tunneling pattern for CH<sub>3</sub>NH<sub>2</sub> in the  $\nu_{11}$  asymmetric CH-stretch fundamental has been previously reported at this meeting. In the experimental pattern, the degenerate species ( $E_1$  and  $E_2$ ) are at the top and bottom of the multiplet and the non-degenerate species (B and A) are between them. In this work, we present two models for the torsion-inversion tunneling behavior in the CH-stretch excited states. Each model includes the lowest order torsion-inversion-vibration interactions available in the context of the model. The first model, which extends Hougen's treatment of methanol,<sup>*b*</sup> couples the two vibrational angular momentum components of the asymmetric CH-stretches to the large-amplitude motion to yield predicted tunneling patterns for the  $\nu_2$  and  $\nu_{11}$  fundamentals. This model gives similar patterns for  $\nu_2$  and  $\nu_{11}$ , in which  $E_1$  and  $E_2$  are in the middle of the multiplet and the non-degenerate species are at the top and bottom. The second model, which follows conceptually Wang and Perry's local mode treatment of methanol,<sup>*c*</sup> couples the three local CH-stretches to each other and to the large-amplitude motion to yield the tunneling patterns for the  $\nu_2$ ,  $\nu_3$  and  $\nu_{11}$  fundamentals. For this model, we found that, for  $\nu_2$  and  $\nu_{11}$ , both  $E_1$  and  $E_2$  are at the bottom of the multiplet, in contrast to  $\nu_3$  and the ground state where they are at the top. The fact that neither model reproduces the observed tunneling pattern for  $\nu_{11}$ , suggests that additional isolated perturbations or systematic interactions are present in the

<sup>a</sup>V.V. Ilyshin et al.J. Mol. Spectrosc. <u>251</u>(56-63), 2008.

<sup>&</sup>lt;sup>b</sup>J.T. Hougen J. Mol. Spectrosc. <u>207</u>(60-65), 2001.

<sup>&</sup>lt;sup>c</sup>X. Wang and D.S. Perry J. Chem. Phys. <u>109</u>(10795-10805), 1998.